

Probing dark energy and inflation with 21cm line observations

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Plan of talk

- Introduction: 21cm line signals from IGM and minihalos
- Applications:
 - ▶ Primordial perturbations
 - ▶ Dark energy
- Summary

Refs:

- ❖ M. Kawasaki, TS, T. Takahashi [1104.5591]
- ❖ K. Kohri, Y. Oyama, T. Sekiguchi & T. Takahashi [arXiv:(1404.4847), (1303.1688), 1608.01601]
- ❖ (H. Tashiro, T. Sekiguchi & J. Silk, N. Sugiyama [arXiv:1311.3295])
- ❖ T. Sekiguchi, T. Takahashi, H. Tashiro & S. Yokoyama [arXiv:1705.00405,1807.02008]

Introduction

CMB last scattering $z \sim 10^3$

Reionization Era?
of the Cosmic History

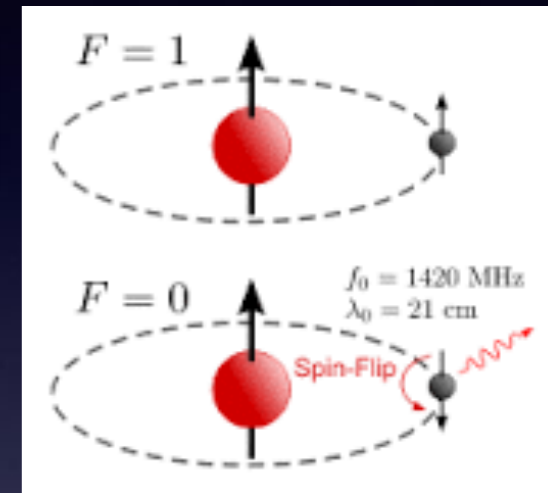
Dark ages (Epoch of reionization)

✓ neutral hydrogen

→ redshifted 21cm line

- tracer of matter fluctuations

- tomography → 3D mapping



reionization $z \sim 6$

~ 300 thousand

~ 500 million

~ 1 billion

~ 9 billion

~ 13 billion

← The Big Bang

← Dark Ages start

Galaxies and Quasars
begin to form
The Reionization starts

The Cosmic Renaissance
The Dark Ages end

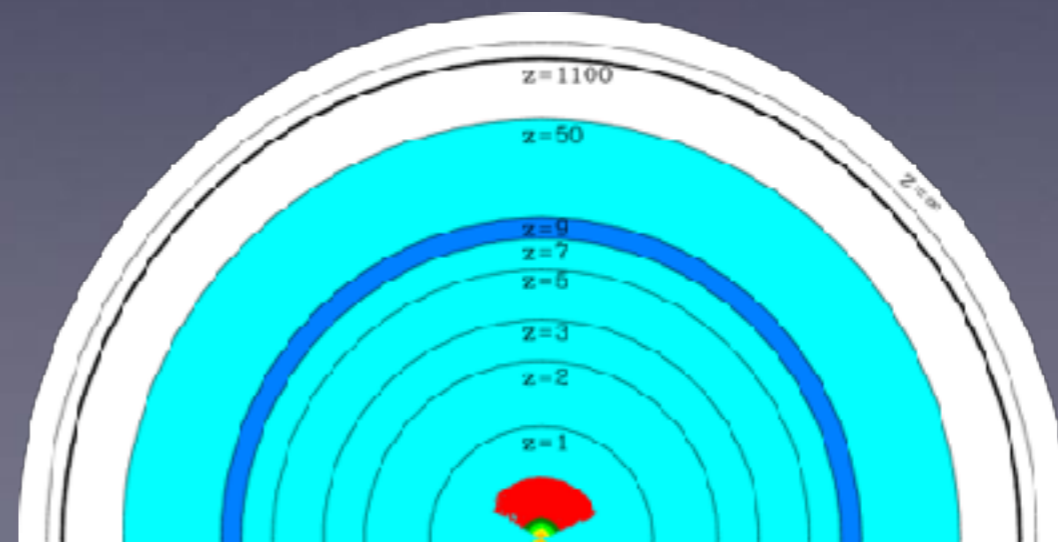
← Reionization complete,
the Universe becomes
transparent again

Galaxies evolve

The Solar System forms

Today: Astronomers
figure it all out!

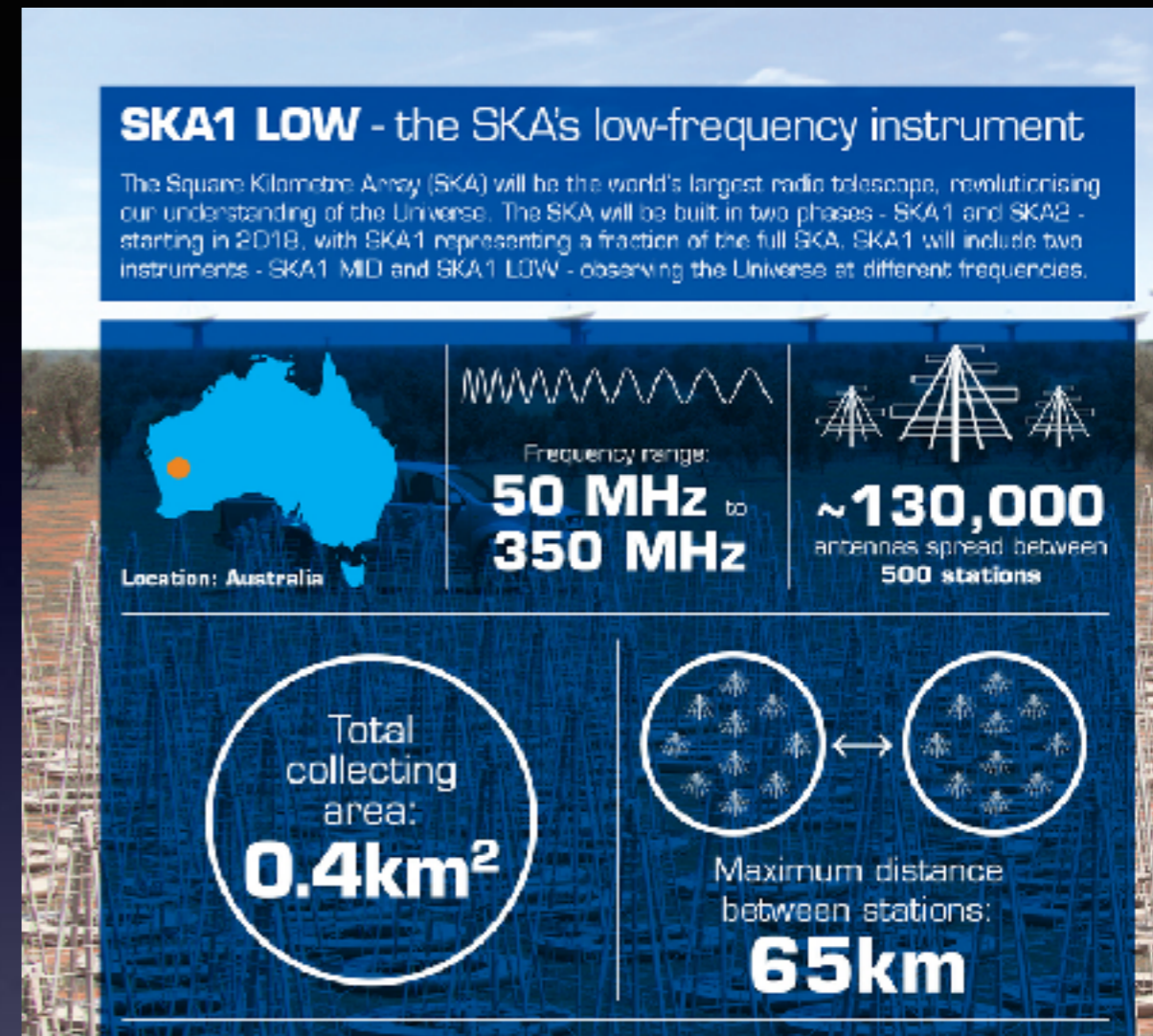
S.G. Djorgovski et al. & Digital Media Center, Caltech



Mao+ '08

Redshifted 21cm line surveys

- Ongoing
 - LOFAR, MWA, ...
- Near future
 - ▶ Square Kilometer Array (SKA-low)
 - 21cm line from $3 < z < 27$;
 - phase1 will start by 2023
 - ▶ Hydrogen Epoch Reionization Array (HERA)
 - main target: 21cm line from $7 < z < 12$
- Far future
 - FFTT, Omniscope, Lunar telescope?



(c) www.skatelescope.org



What do 21 cm surveys observe?

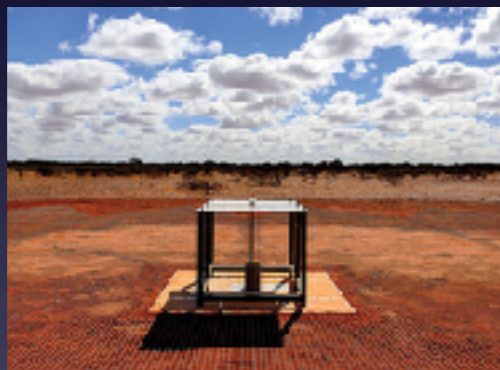
- Spin temperature

- Ratio of triplets to singlet $n_{\text{triplet}}/n_{\text{singlet}} = 3 \exp[-E_{21\text{cm}}/T_s]$

What do 21cm surveys observe?

- Spin temperature

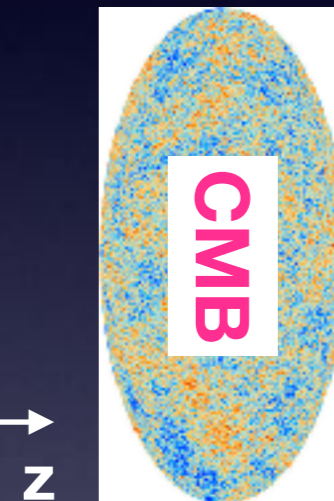
- Ratio of triplets to singlet $n_{\text{triplet}}/n_{\text{singlet}} = 3 \exp[-E_{21\text{cm}}/T_s]$



neutral hydrogen



$$1 + z_\nu = 1.4\text{GHz}/\nu$$



- Brightness temperature

- Radiative transfer

$$T_{21\text{cm}}(\nu) = \frac{T_s(z_\nu) - T_{\text{CMB}}(z_\nu)}{1 + z_\nu} (1 - e^{-\tau_{21\text{cm}}(\nu)}) \simeq (T_s - T_{\text{CMB}})\tau_{21\text{cm}}$$

Emission if $T_s > T_{\text{CMB}}$
Absorption if $T_s < T_{\text{CMB}}$

- 21cm optical depth

$$\tau_{21\text{cm}}(\nu) = \int dl \frac{3A_{10}\lambda_{21\text{cm}}^2}{32\pi} \frac{n_{\text{HI}}(z_\nu)}{T_s(z_\nu)} \phi(\nu)$$

Sources of 21cm line



21 cm fluctuation from IGM

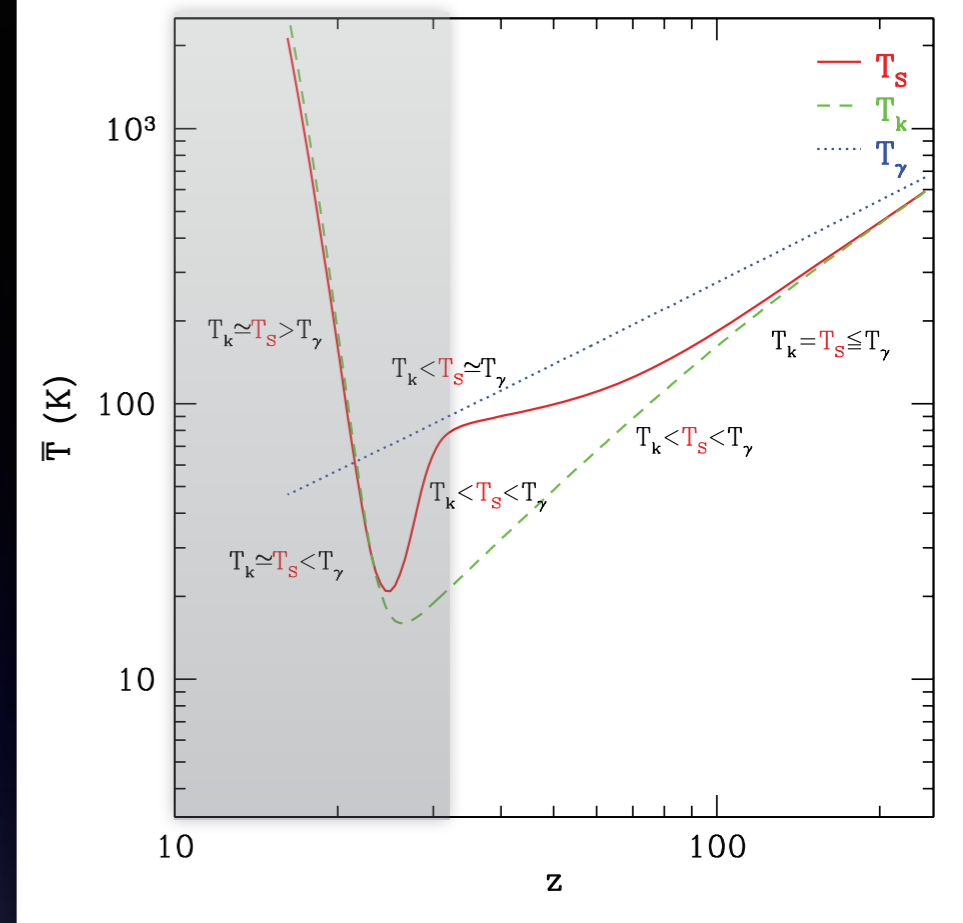
- Differential brightness temperature

$$\Delta T_b = T_b - T_{\text{CMB}} \simeq \frac{T_s - T_{\text{CMB}}}{1+z} \tau_{21\text{cm}}$$

At high redshifts $z > 20$

(prior to formation of first objects), $T_s < T_{\text{CMB}}$.

→ absorption



Messinger et al. 2011

- Fluctuations in 21 cm brightness temperature

$$\delta_{21\text{cm}} \approx \underbrace{\frac{\bar{T}_{\text{CMB}}}{\bar{T}_s - \bar{T}_{\text{CMB}}} (\delta_{T_s} - \delta_{T_{\text{CMB}}}) + \delta_{n_{\text{HI}}}}_{\text{depends on } \delta_b} - \underbrace{\frac{\hat{n} \cdot d\vec{v}_b / dr}{H}}_{\text{depends only on } \delta_m}$$

$$\mu = \hat{k} \cdot \hat{n}$$

\hat{n} : line-of-sight direction

isotropic

$\propto \mu^2$

21 cm can probe δ_b , separately from δ_m .

IGM 21cm power spectrum

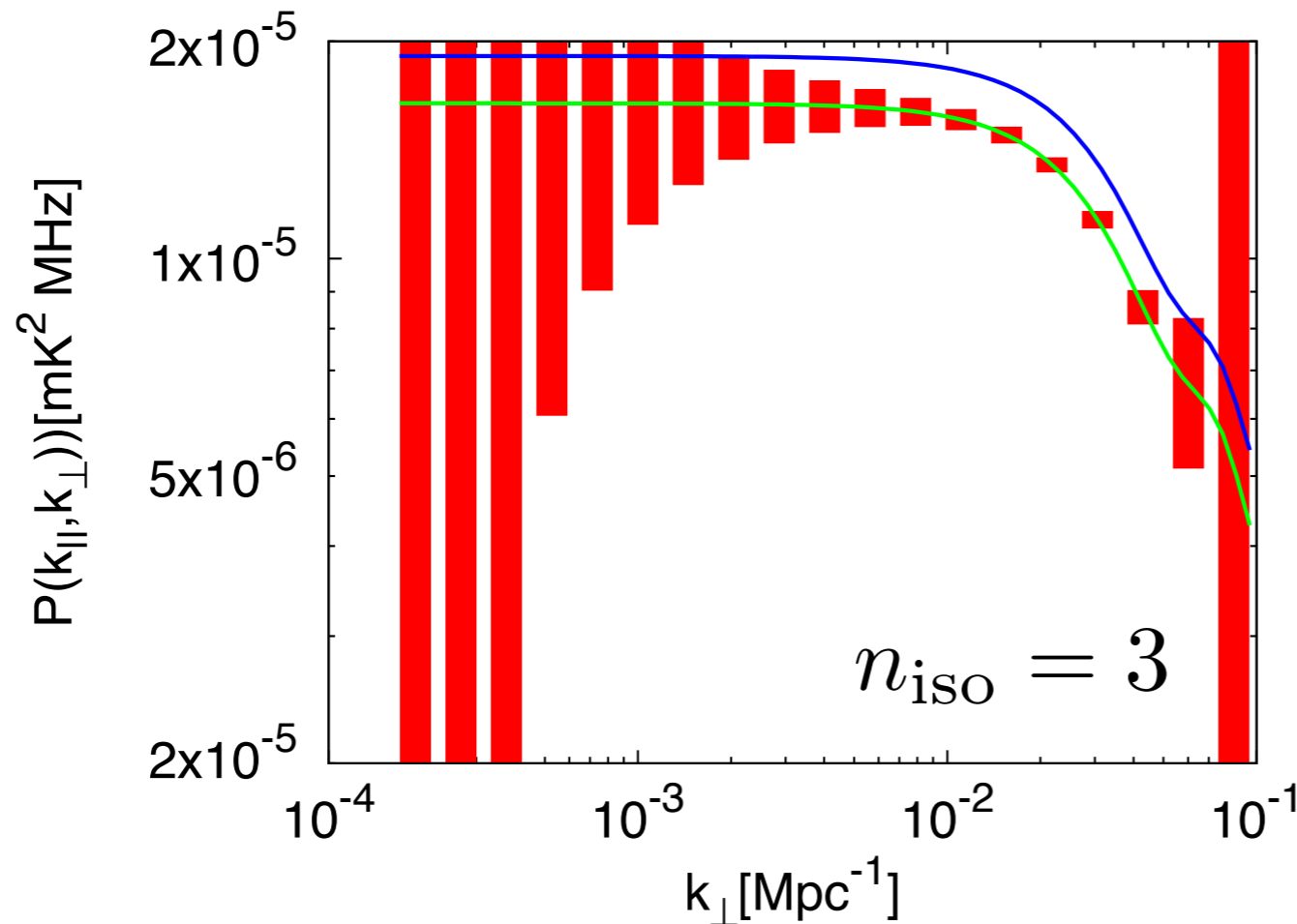
- Tomographic power spectrum

$$\langle \delta T_b(\vec{k}) \delta T_b(\vec{k}') \rangle \equiv (2\pi)^3 \delta(\vec{k}_1 + \vec{k}_2) P(k_{\parallel}, k_{\perp})$$

$$P(k_{\parallel}, k_{\perp}) = P_{\delta\delta}(k) + 2\mu^2 P_{\delta v}(k) + \mu^4 P_{vv}(k) \quad k = \sqrt{k_{\parallel}^2 + k_{\perp}^2} \quad \mu = k_{\parallel}/k$$

- $P_{S_m}(k_0)/P_{AD}(k_0) = 0.1$ at $k_0 = 0.0002 \text{Mpc}^{-1}$

$z = 50, k_{\parallel} = 0.06 \text{Mpc}^{-1}$



pure CI — pure BI
 FFTT sensitivity

CDM/baryon isocurvature perturbation can be distinguished by 21cm.


Distinguishing CDM and baryon isocurvature


Kawasaki, TS, Takahashi 2011

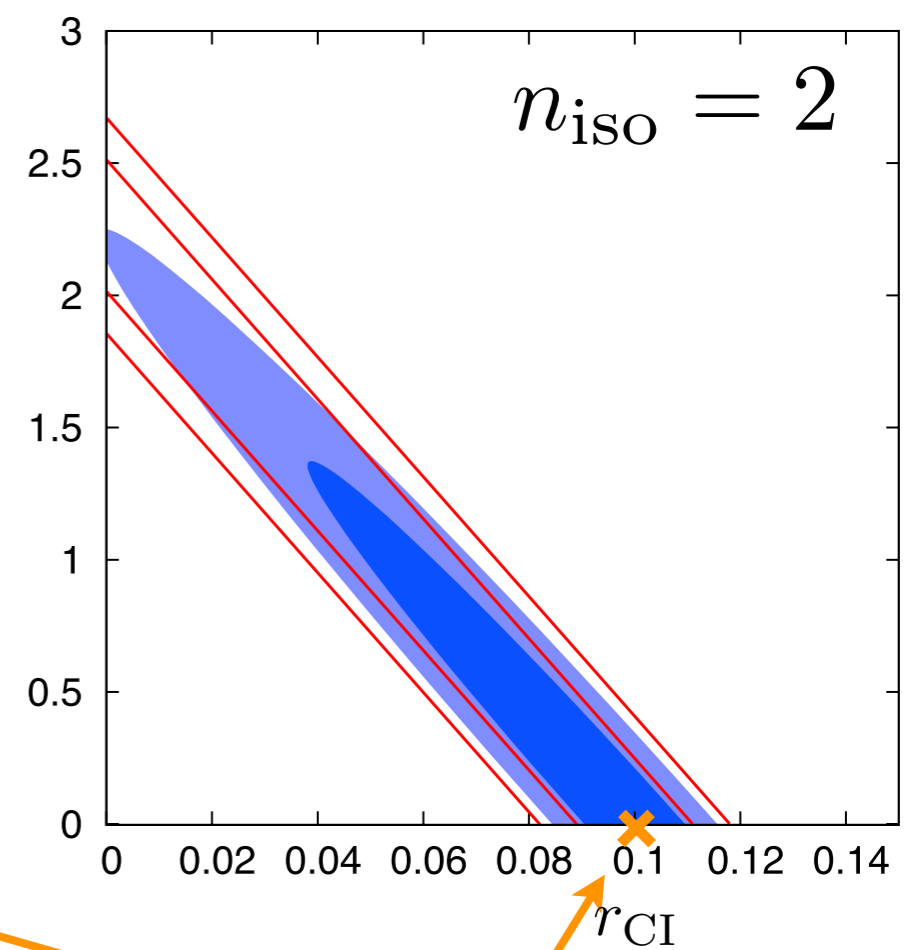
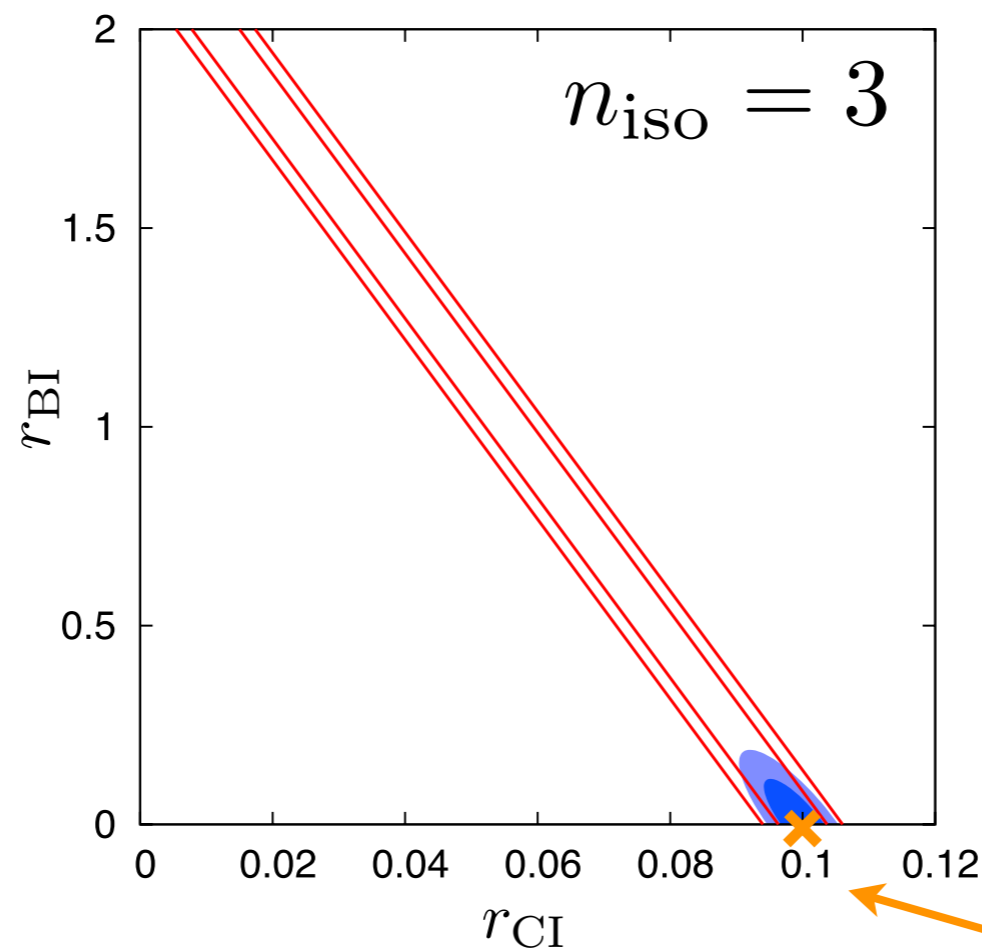
- Fisher matrix analysis

$$r_{\text{CI}} = P_{\text{CI}}(k_0)/P_{\text{adi}}(k_0), \quad r_{\text{BI}} = P_{\text{BI}}(k_0)/P_{\text{adi}}(k_0) \quad k_0 = 0.002\text{Mpc}^{-1}$$

2d constraints

CMBpol alone 

CMBpol+FFTT
@ z=30, 40, 50 



*fiducial model:
pure CI
 $(r_{\text{CI}}, r_{\text{BI}}) = (0.1, 0)$

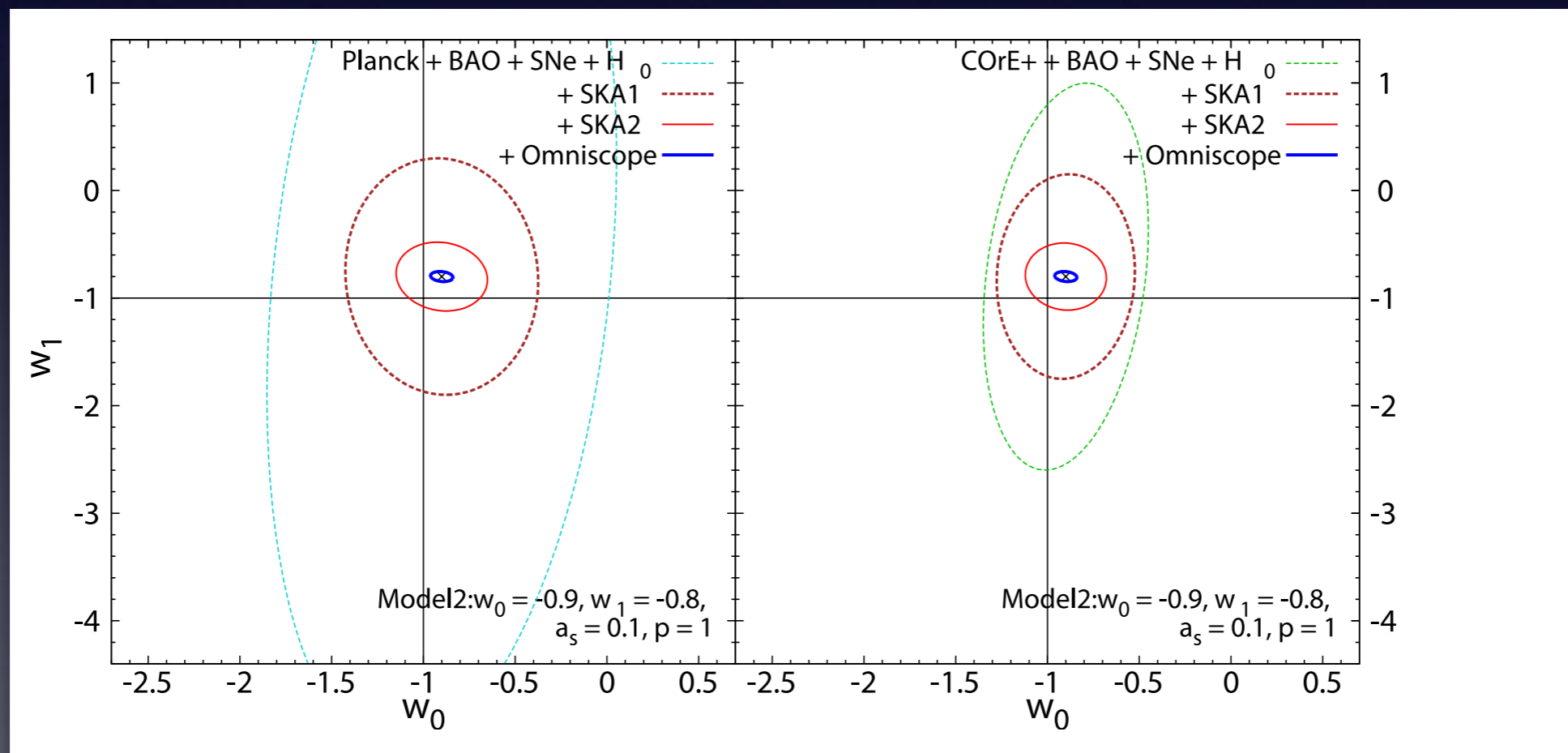
Future 21 cm surveys can distinguish CI/BI if $n_{\text{iso}} \gtrsim 2$.

Dark energy

Kohri, Oyama, TS, Takahashi 2017

Parameterized EoS

$$w(z) = w_0 w_1 \frac{a^p + a_s^p}{w_1 a^p + w_0 a_s^p} \approx \begin{cases} w_0 & (\text{for } a \gg a_s) \\ w_1 & (\text{for } a \ll a_s) \end{cases}$$



Redshifted 21cm line fluctuations can constrain early-type dark energy better than CMB.

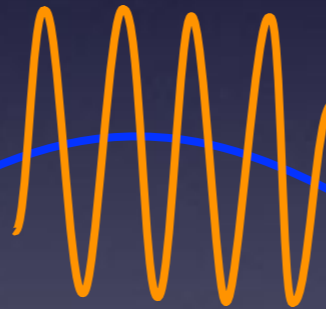
Sources of 21cm line

✓ Minihalos

Iliev+ '02;
Furlanetto & Loeb '02

δ_b

✓ Smooth IGM



Minihalos

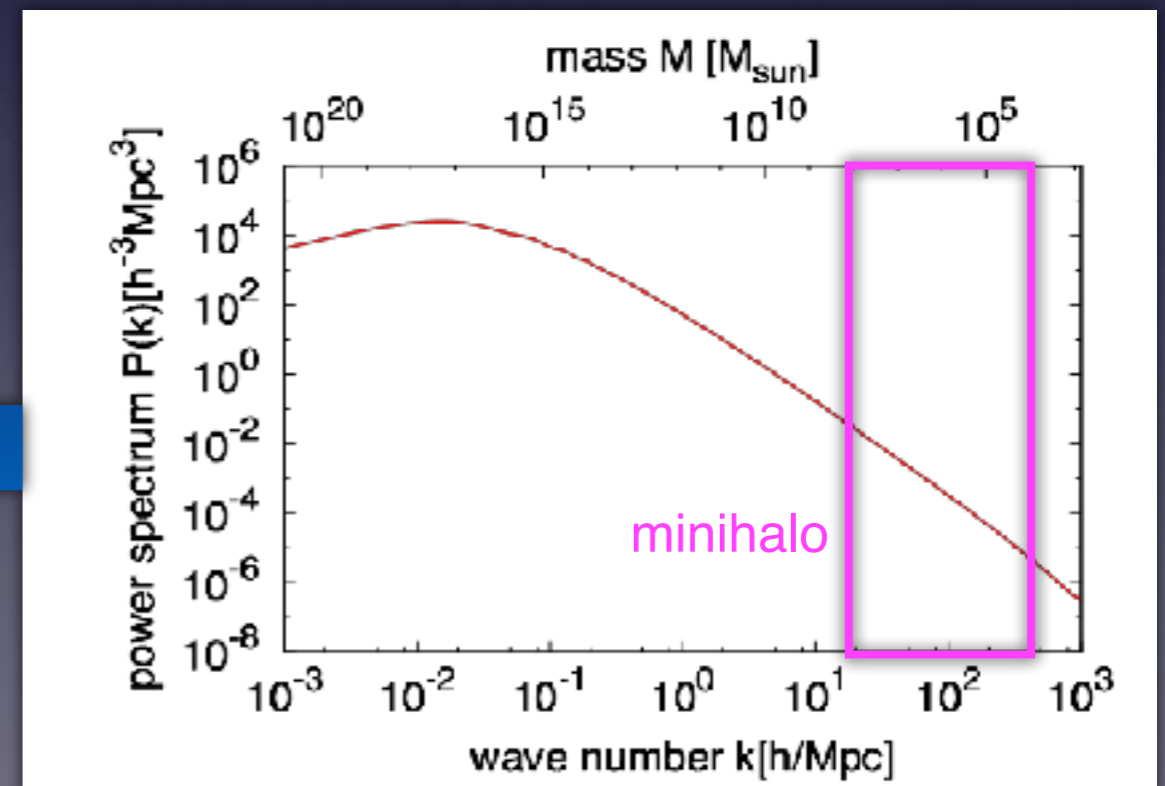
Halos too small to host galaxies

- No star formation: $T_{\text{gas}} < 10^4 \text{K}$ (inefficient radiative cooling)
 - dense neutral hydrogen inside; resistant to ionization

- Mass: $10^4 M_{\text{sun}} < M < 10^8 M_{\text{sun}}$

Sensitive to small-scale
($< 0.1 \text{Mpc}$) fluctuations

- Abundant, even at high- z

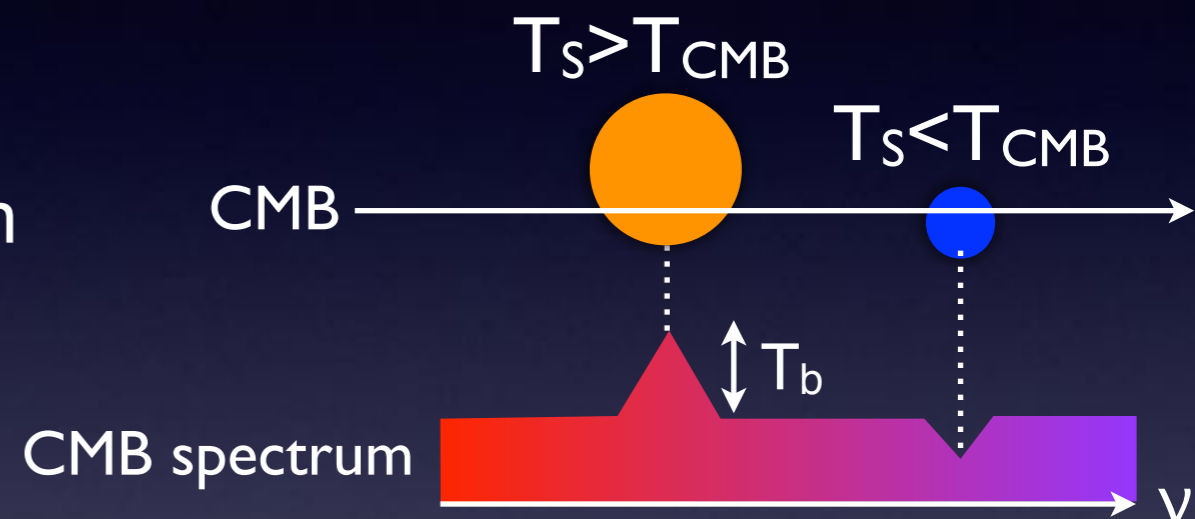


21cm line signal from minihalos

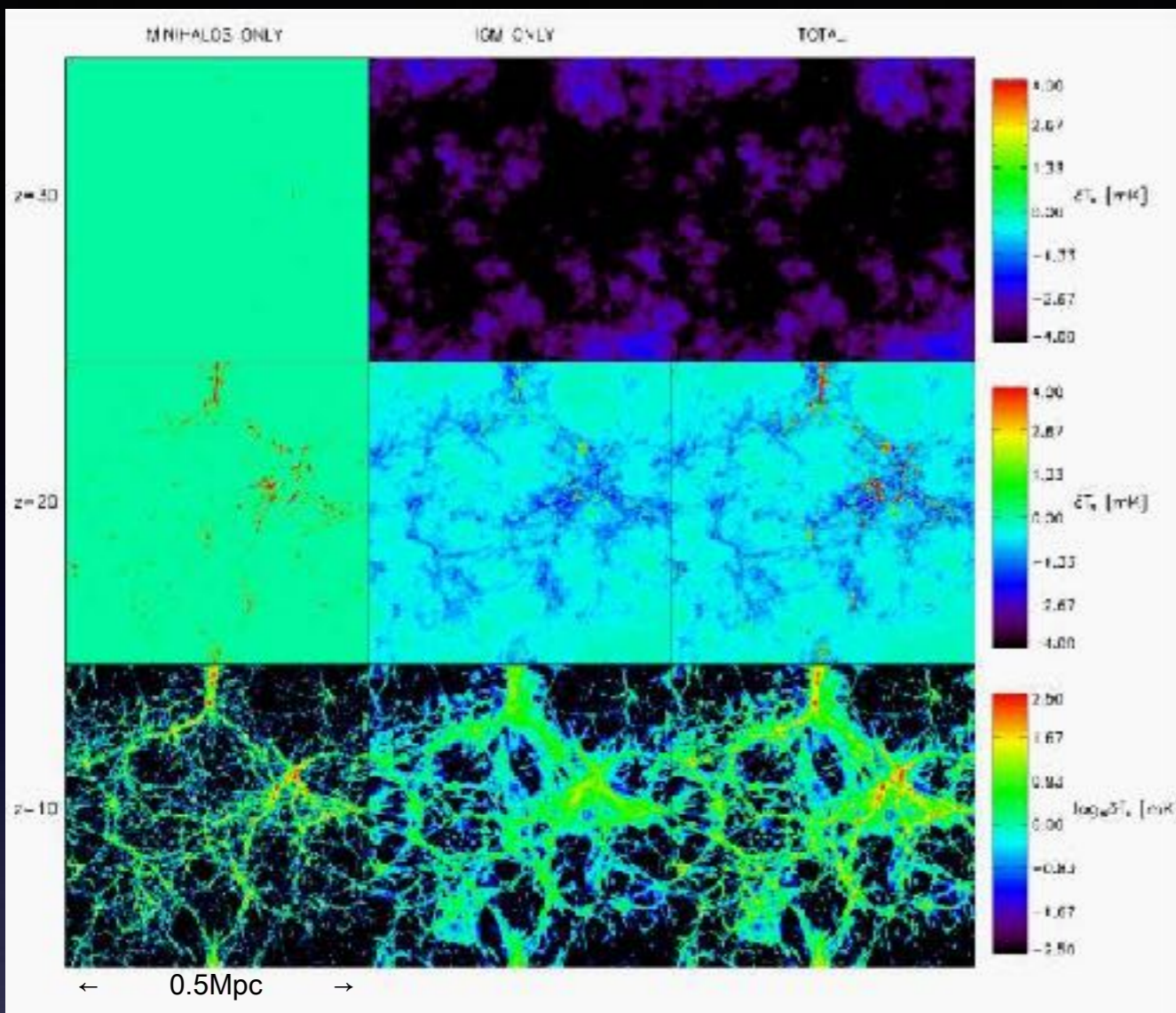
Iliev+ '02;
Furlanetto & Loeb '02

“21cm forest” in CMB

- Minihalos create emission/absorption features in CMB spectrum at radio frequency $\nu = 1.4\text{GHz}/(1+z)$
- Large (small) halos appear as emission (absorption)
- Individual halos are too small (size $\sim \text{kpc}$) to be resolved
→ intensity maps (like CMB)



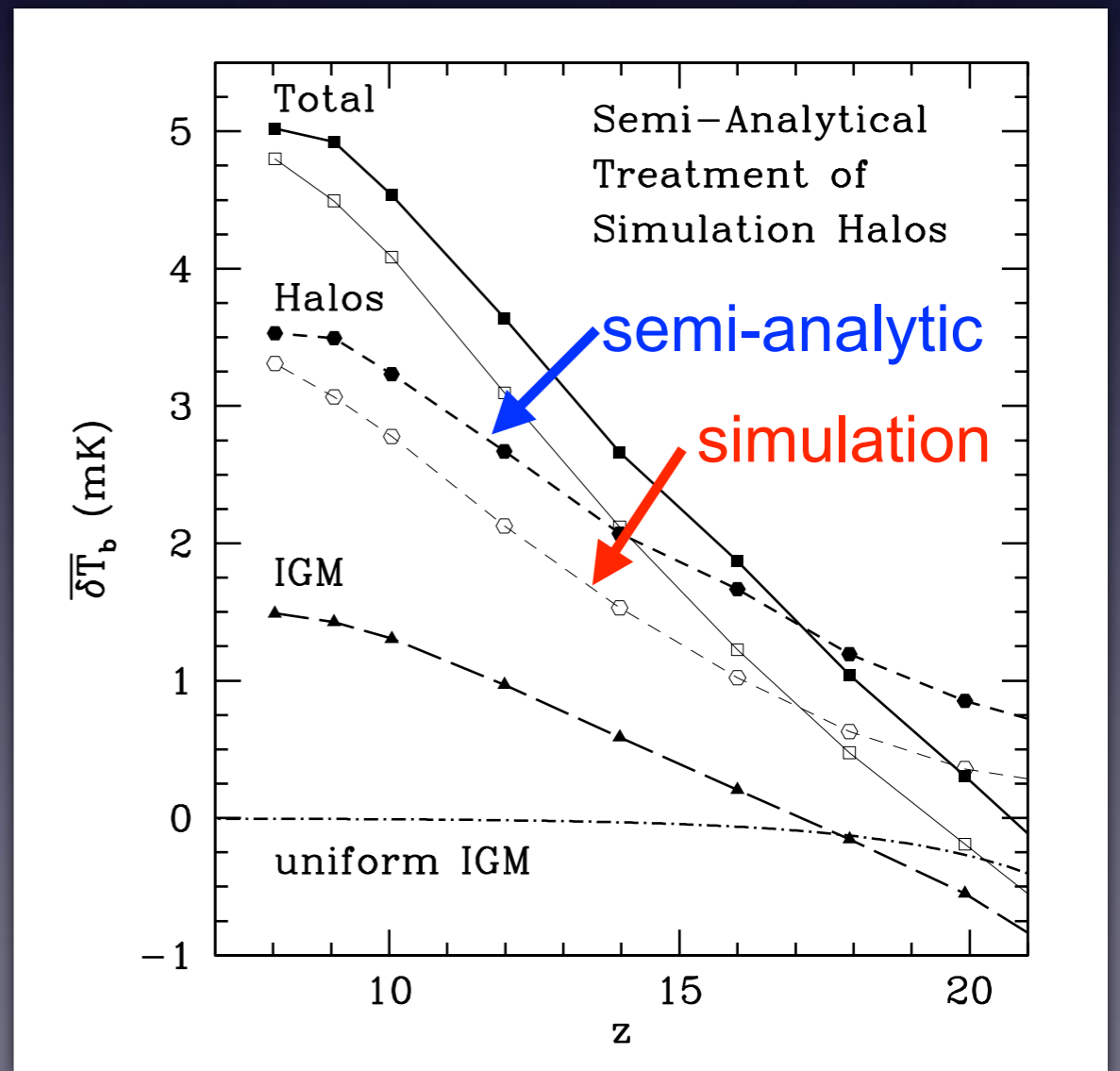
Simulations



N-body+hydro simulations Shapiro+ '06

Minihalos can exceed the IGM around the epoch of reionization

Semi-analytical description agrees with simulations



21cm angular power spectrum from minihalos (1)

Iliev+ '02; TS, Takahashi, Tashiro & Yokoyama '17

Tomographic anisotropy (w/o redshift space distortion)

$$\delta T_b(\hat{n}, \nu) = \left[\int_{M_{\min}}^{M_{\max}} dM \underbrace{T_b^{(\text{single})}(M, z_\nu)}_{\text{strength of 21cm line emission/absorption from single minihalo}} \underbrace{\frac{dN(M, z_\nu)}{dM}}_{\text{mass function}} \underbrace{b(M, z)}_{\text{halo bias}} \right] \delta(\vec{x} = r_\nu \hat{n}, z_\nu)$$

sensitive to small-scale (<Mpc) matter fluctuations
matter fluctuations at large scales (>Mpc)

21cm angular power spectrum from minihalos (2)

Iliev+ '02; TS, Takahashi, Tashiro & Yokoyama '17

Redshift-space distortion (Kaiser effect)

$$\delta T_b(\hat{n}, \nu) = \bar{T}_b(z) [\beta(z) + f(z)\mu^2] \delta(\vec{x}, z) \quad \text{with } \mu = \hat{k} \cdot \hat{n}$$

mean signal: $\bar{T}_b(z) = \int dM \mathcal{F}(M, z) = \int dM T_b^{(\text{single})}(M, z) \frac{dN}{dM}(M, z)$

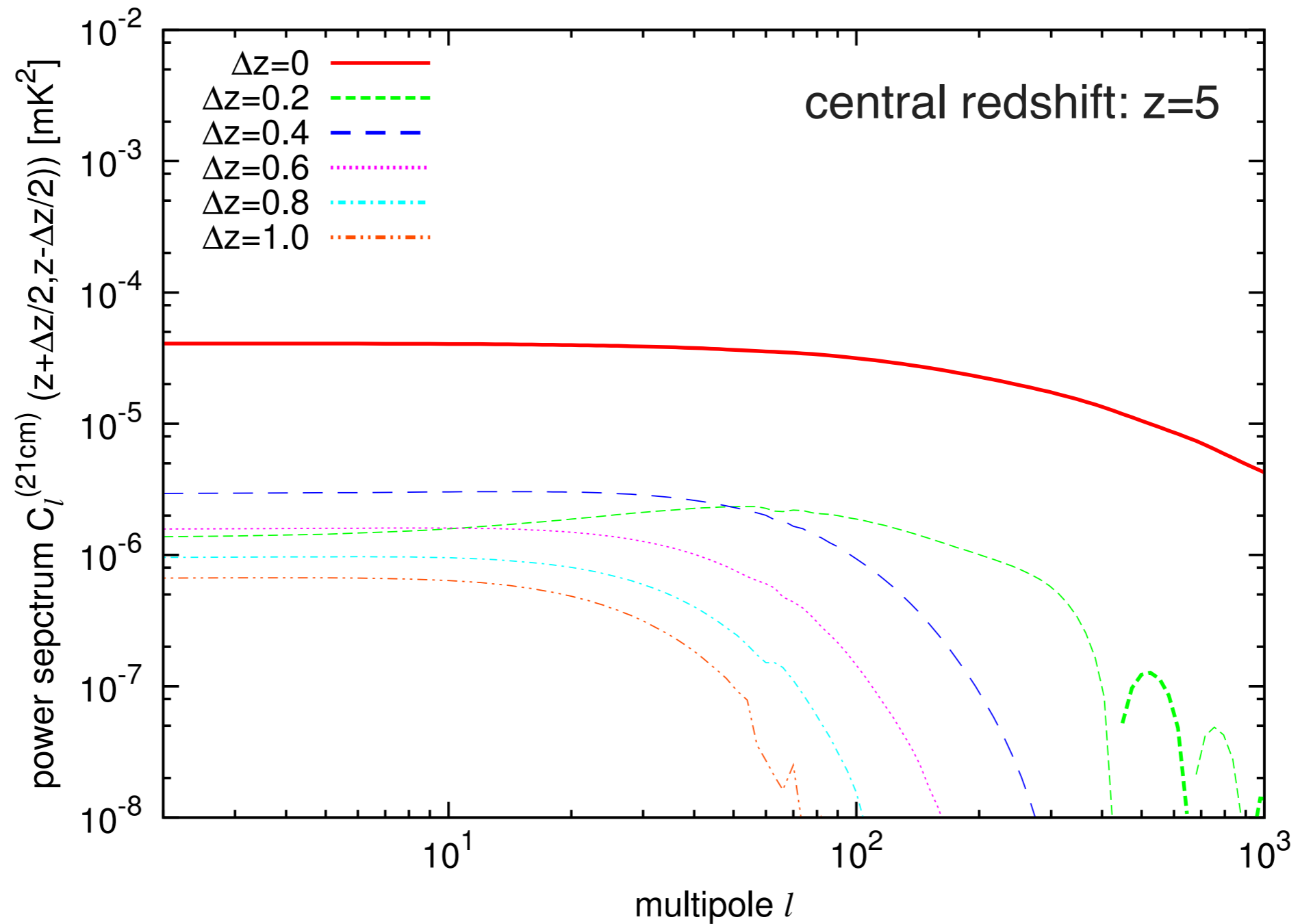
growth rate: $f(z) = d \ln D(z) / d \ln a$

flux-weighted effective bias: $\beta(z) = \frac{1}{\bar{T}_b(z)} \int dM \mathcal{F}(M, z) b(M, z)$

Tomographic angular power spectrum

$$C_l(z, z') = \frac{1}{2l+1} \sum_m a_{lm}(z) a_{lm}^*(z') \quad \text{with } a_{lm}(z_\nu) = \int d\hat{n} \delta T_b(\hat{n}, \nu) Y_{lm}^*(\hat{n})$$

21cm angular power spectrum from minihalos (3)



Application (1): Primordial spectral runnings

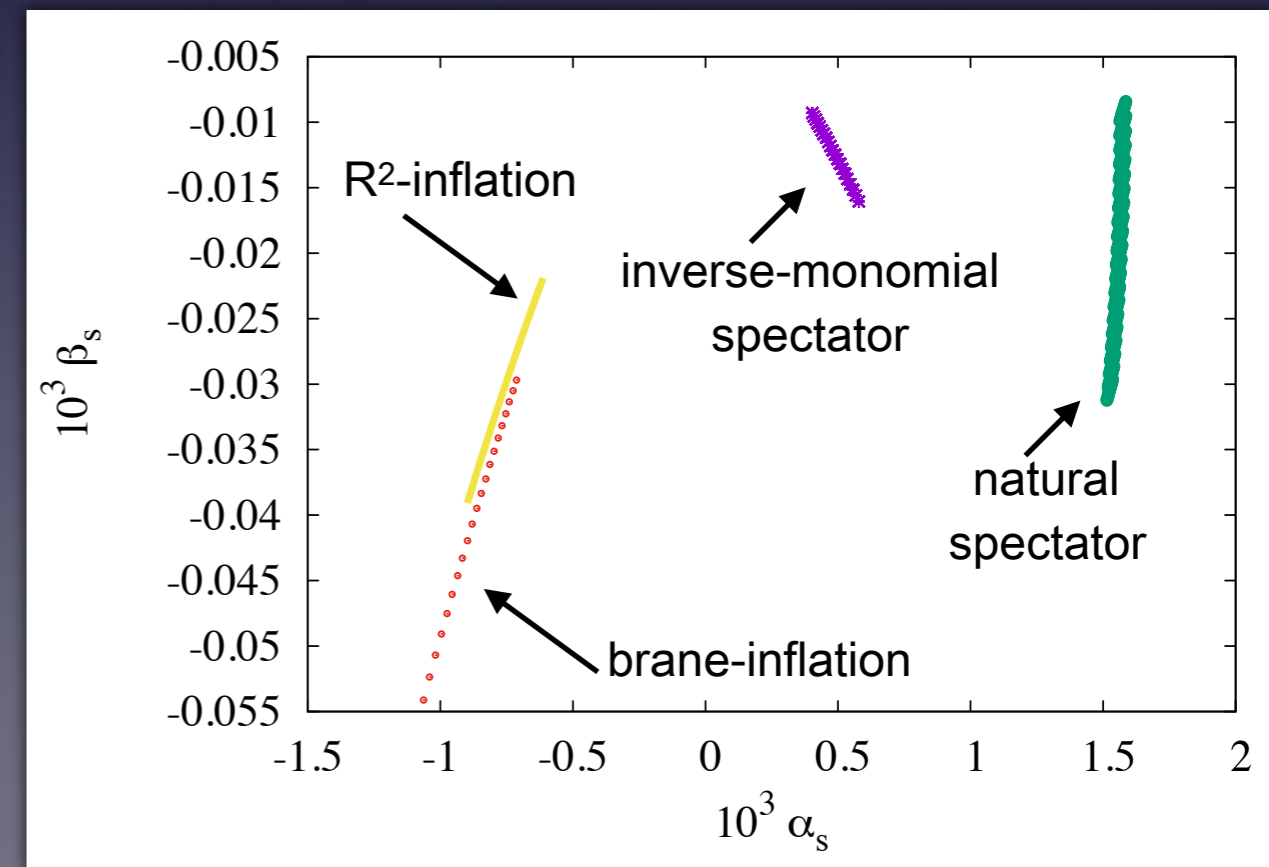
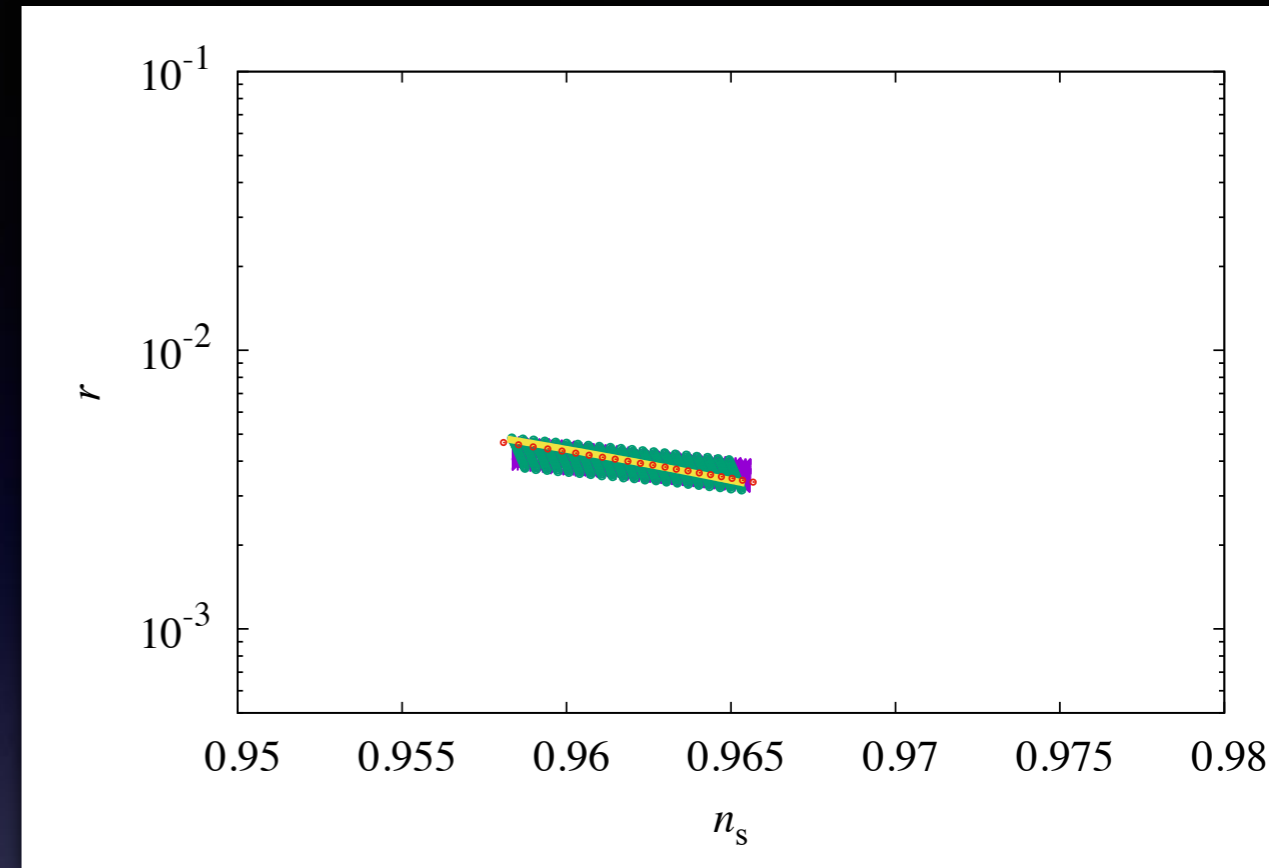
TS, Takahashi, Tashiro & Yokoyama [arXiv:1705.00405]

Spectrum of primordial fluctuations

$$\mathcal{P}(k) \propto k^{n_s - 1 + \frac{1}{2}\alpha_s \ln(k/k_*) + \frac{1}{6}\beta_s \ln^2(k/k_*) + \dots}$$

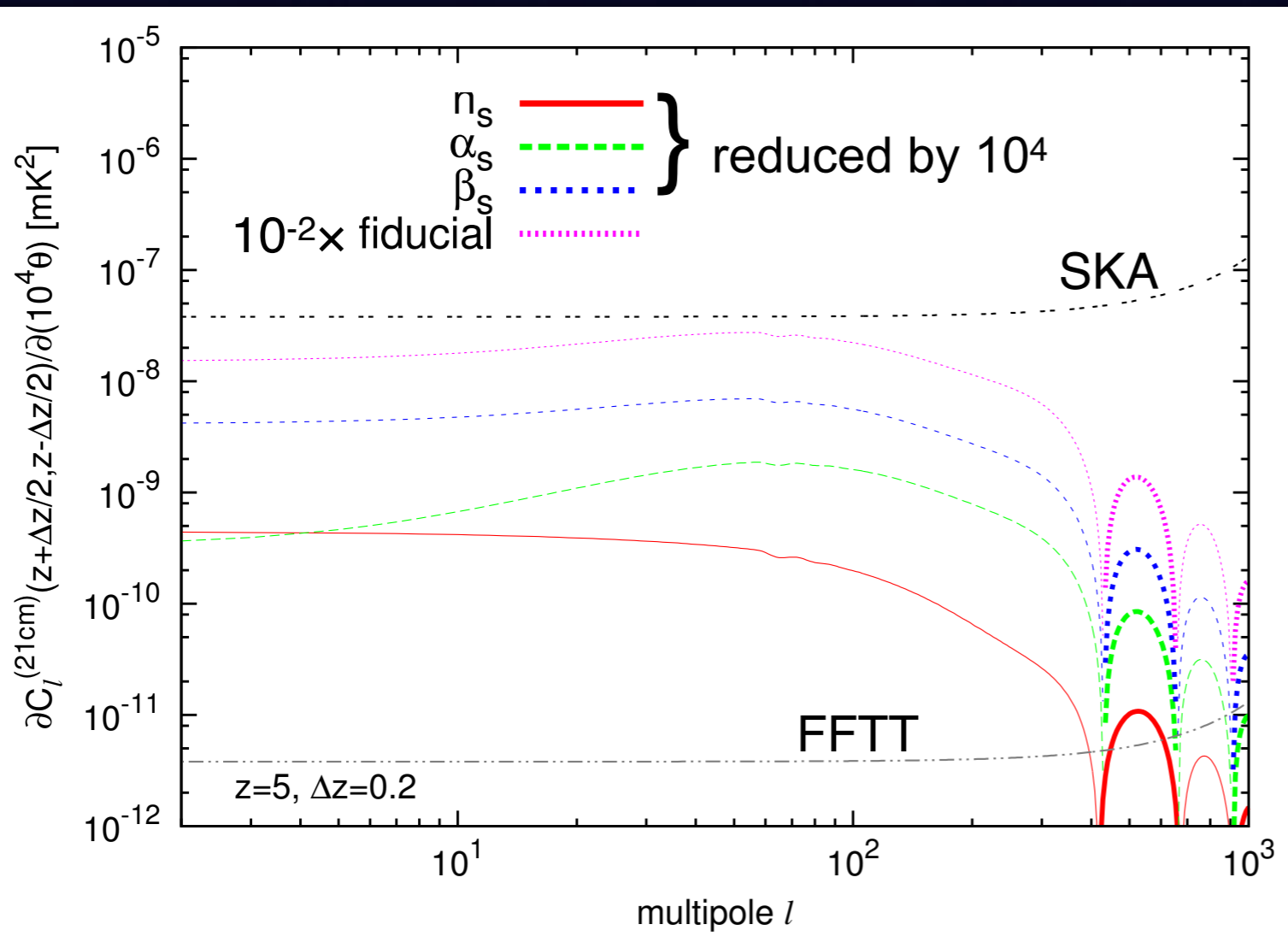
- Many models degenerate in the n_s - r plane
- However, they can be distinguished from the scale dependence of n_s

Spectral runnings: a key observable for discriminating inflation models



Application (1): Primordial spectral runnings (cont'd)

Parameter response



- Lower order spectral parameters (e.g. n_s or α_s) → spectral shapes
- Higher order parameters (e.g. β_s) → overall amplitudes
- Solves parameter degeneracy
- Radial scale-dependence also enhances the discrimination

Application (1): Primordial spectral runnings (cont'd)

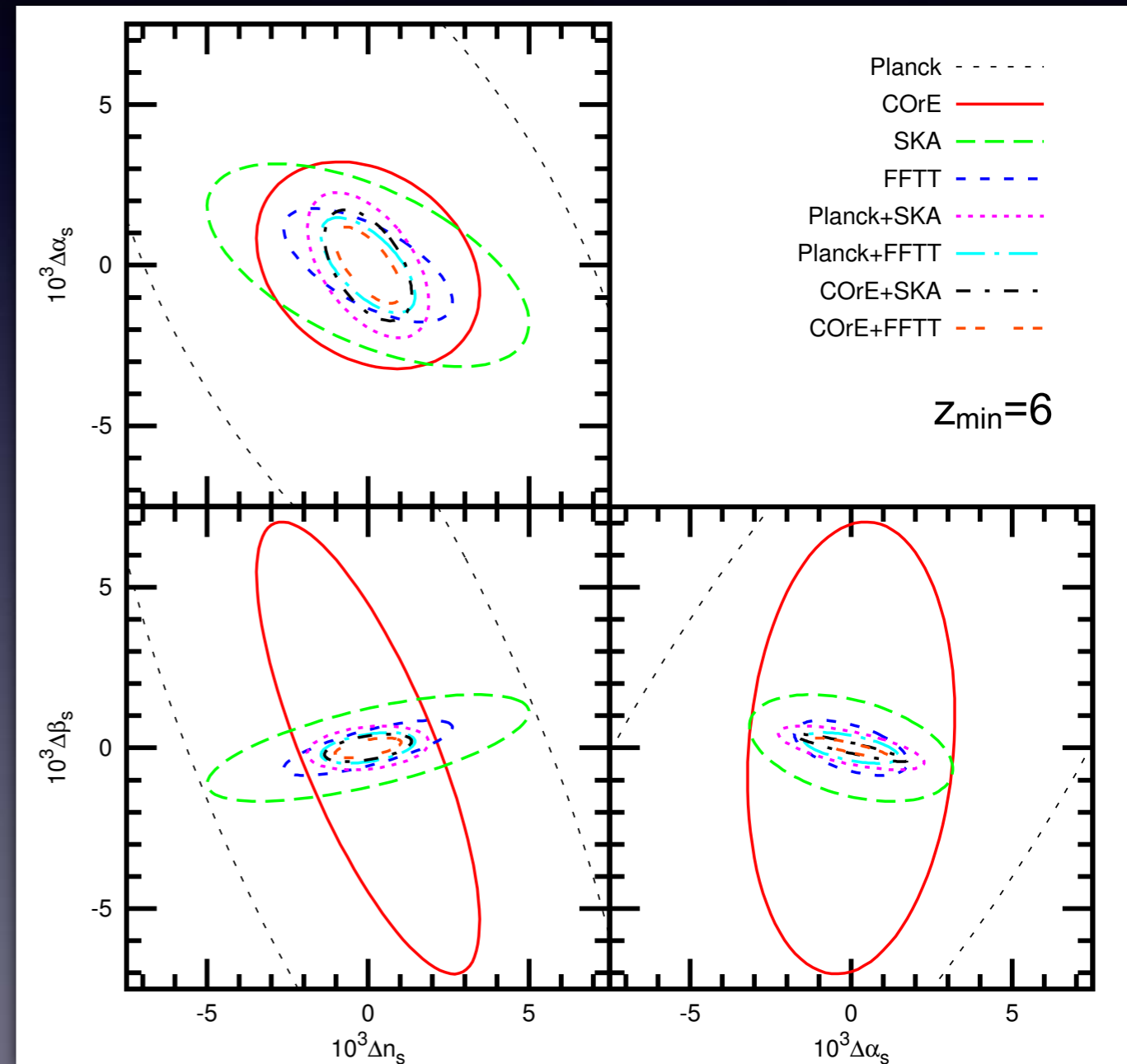
Forecasted constraints

Combination of CMB and 21cm is beneficial due to lever-arm effect.

$$\Delta\alpha_s=10^{-3}, \Delta\beta_s=10^{-4}$$

Constraints are dependent on z_{\min} only mildly.

	z_{\min}	$10^{-3}\Delta n_s$	$10^{-3}\Delta\alpha_s$	$10^{-3}\Delta\beta_s$
Planck+SKA	4	1.4	1.4	0.40
	6	1.7	2.0	0.63
	8	2.3	3.0	0.85
	10	3.6	4.7	1.2
COre+FFTT	4	0.85	0.96	0.24
	6	0.95	1.1	0.28
	8	1.0	1.2	0.31
	10	1.1	1.3	0.33



Application (2): Primordial non-Gaussianity

TS, Takahashi, Tashiro & Yokoyama, in prep.

Local type non-Gaussianity:

$$\Phi(\vec{x}) = \Phi_G(\vec{x}) + f_{\text{NL}}(\Phi_G(\vec{x})^2 - \langle \Phi_G \rangle^2) + g_{\text{NL}}\Phi_G(\vec{x})^3$$

- Small in single field inflation: $f_{\text{NL}} \sim \mathcal{O}(0.01)$, $g_{\text{NL}} < \mathcal{O}(10^{-3})$
- Large in multi-field models (e.g. curvaton, modulated reheating, etc.)

Current tightest bound (Planck 2015)

$$f_{\text{NL}} = 0.8 \pm 5.0, \quad g_{\text{NL}} = (9.0 \pm 7.7) \times 10^4$$

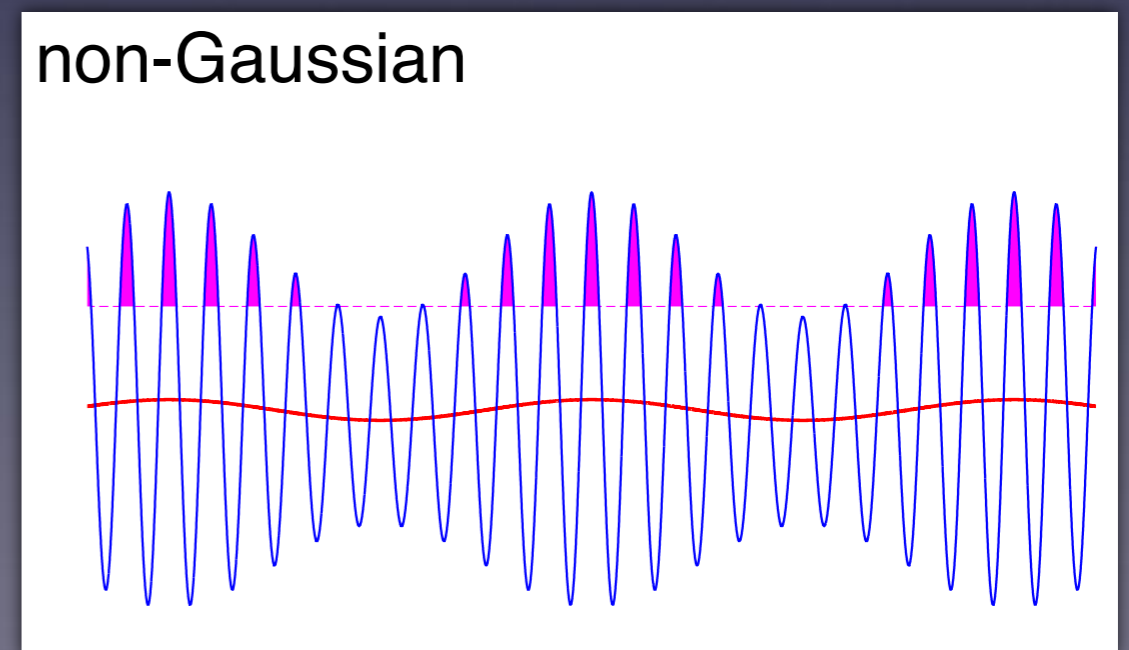
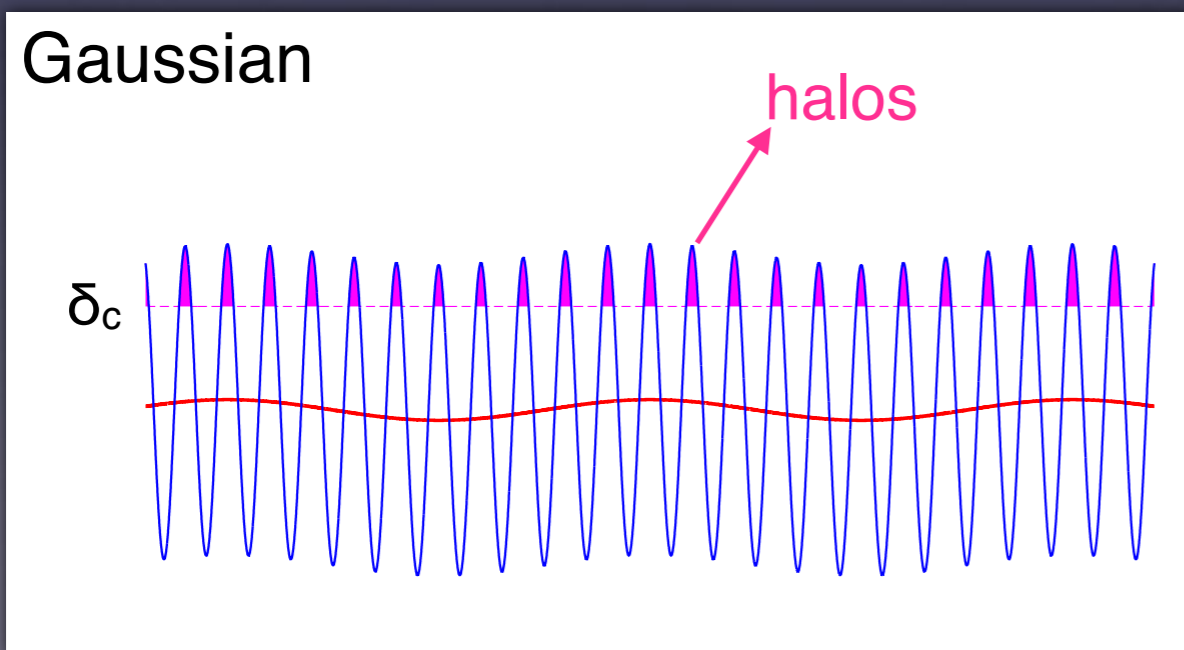
cf. $g_{\text{NL}} = (-3.3 \pm 2.2) \times 10^5$ (WMAP 9yr)

TS & Sugiyama '13

Application (2): Primordial non-Gaussianity (cont'd)

Effects of local-type non-Gaussianity on (mini)halos

- Correlation between large and small scale fluctuations
 - relative halo # count $\frac{n_{\text{halo}}}{\rho_m}(\vec{x})$ is modulated by large-scale fluctuations
- **scale-dependent halo bias** Dalal+ '08; Slosar+ '08



Application (2): Primordial non-Gaussianity (cont'd)

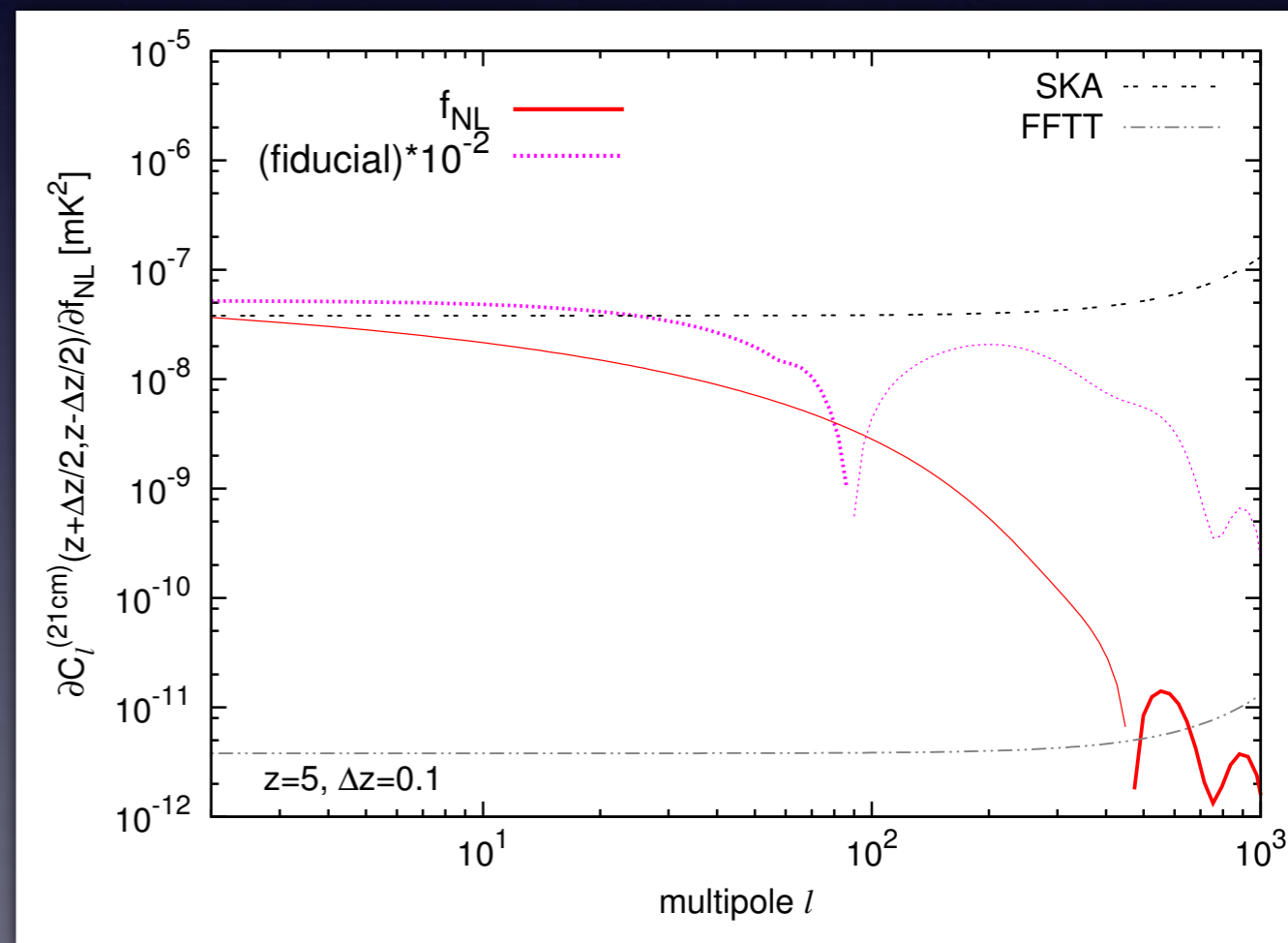
Effects on minihalo power spectrum

Bias is more enhanced at larger scales

$$\Delta\beta(k, z) \approx \{\beta_f(z)f_{\text{NL}} + \beta_g(z)g_{\text{NL}}\} \frac{3\Omega_m H_0^2}{2k^2 T(k) D(z)}$$

→ 21cm line surveys are advantageous

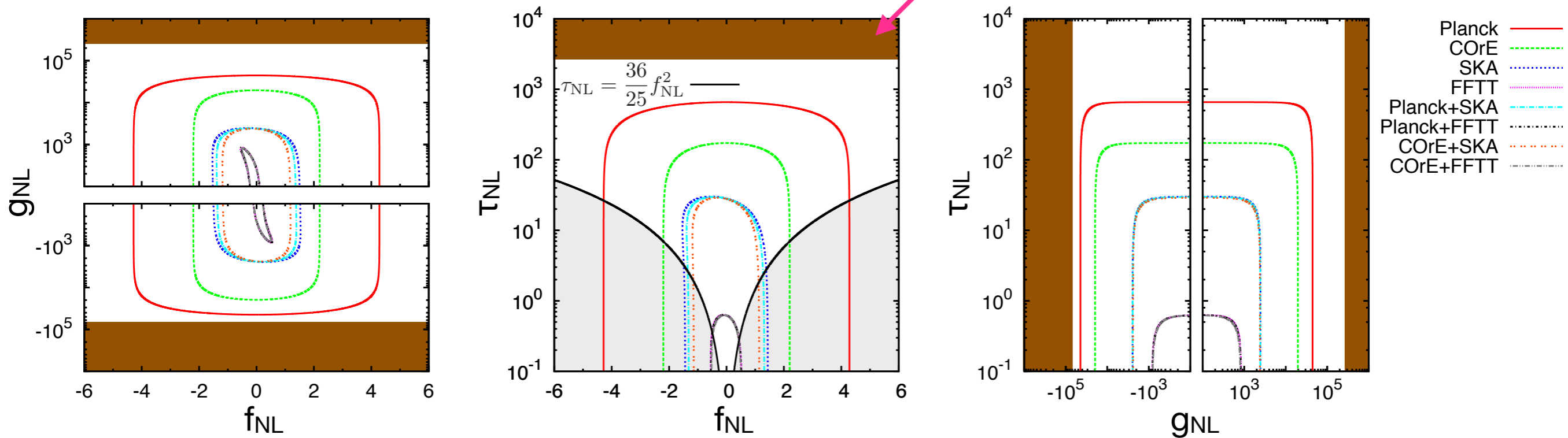
- ✓ large transverse scale comparable to CMB
- ✓ cross-correlation of different redshifts



Application (2): Primordial non-Gaussianity (cont'd)

Forecasted constraints

Current CMB bound



- Minihalos can improve the current (CMB) bound by orders of magnitude

$$\Delta g_{\text{NL}} \simeq O(10^3), \Delta \tau_{\text{NL}} \simeq O(10) \quad (\text{for SKA})$$

- Suyama-Yamaguchi inequality can be tested

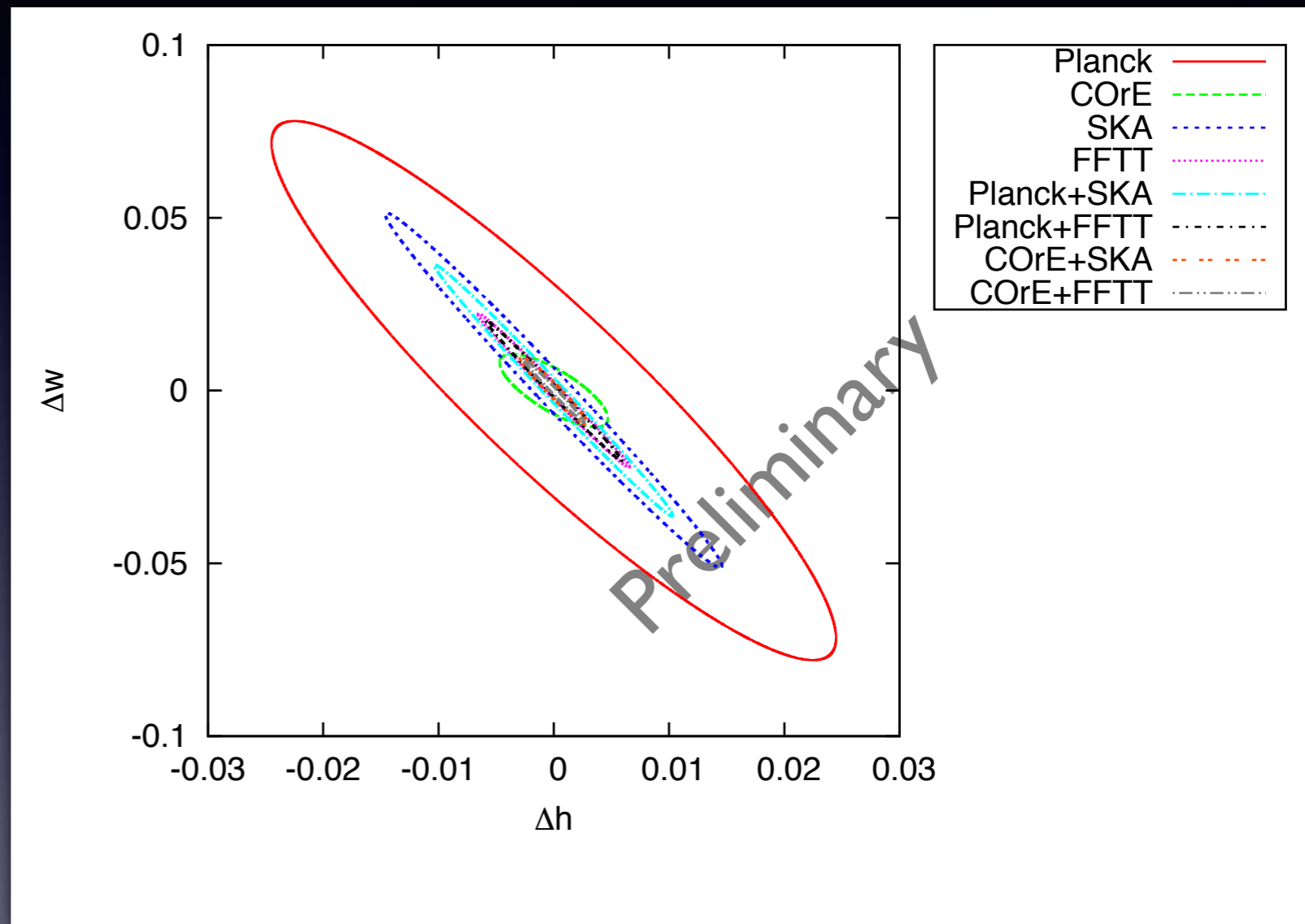
Application to dark energy

Constant EoS

- SKA: $\Delta w \sim 0.05$
- FFTT: $\Delta w = 0.02$
- Cf. Planck: $\Delta w = 0.08$

Both CMB and 21cm suffers from the degeneracy between w and the Hubble parameter.

Incorporation of direct Hubble measurements may be useful.



We will pursue our analysis with early-type DE in the future.

Summary

- High redshifted 21cm line fluctuations are a novel probe of the cosmological structure. There are largely two types of sources: smooth IGM and minihalos.
- Exploiting the tomographic nature of redshifted 21cm line fluctuations, we can constrain a variety of cosmological models.
 - Primordial fluctuations (spectral runnings, non-Gaussianity, etc.)
 - Dark energy
 - (DM, neutrinos, etc)