## Homework 1

Due 20/03/2024

## 1. Relativity Gymnastics

(a) Show that the invariance of the interval between two events  $(t^2 - |\vec{x}|^2)$  under Lorentz transformations is satisfied if we define the position four-vector  $x^{\mu} \equiv (t; \vec{x})$  such that

$$x^{\prime\mu} = \Lambda^{\mu}_{\nu} x^{\nu}$$

where we defined the Lorentz transformations as  $\Lambda^{\mu}_{\nu}$ . The inverse Lorentz transformation is defined by  $x^{\nu} = \Lambda^{\nu}_{\mu} x^{\prime \mu}$ , and it satisfies

$$\Lambda^{\mu}_{\nu}\,\Lambda^{\nu}_{\rho}=\delta^{\mu}_{\rho}$$

(b) The Minkowski space metric is given by

$$g^{\mu\nu} = g_{\mu\nu} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \ ,$$

- i. A contra-variant four-vector is denoted as  $v^{\mu} = (t; \vec{v})$  Show that if we define a covariant four-vector as  $v_{\mu} \equiv (t; -\vec{v})$ , we can write  $v.v = v^{\mu}v_{\mu}$ .
- ii. Verify that  $v_{\mu} = g_{\mu\nu} v^{\nu}$ , and  $v^{\mu} = g^{\mu\nu} v_{\nu}$ .
- iii. Verify that the interval is given by the position four-vector squared defined by  $x.x \equiv x^{\mu}g_{\mu\nu}x^{\nu}$ .
- (c) Show that the relativistic dispersion relation  $E^2 |\vec{p}|^2 = m^2$ , where *m* is the particle mass, is consistent with defining the momentum four-vector as  $P^{\mu} = (E; \vec{p})$ .

(d) Define the differential operators

$$\frac{\partial}{\partial x^{\mu}} \equiv \partial_{\mu} \equiv (\frac{\partial}{\partial t}; \vec{\nabla}) \qquad \qquad \frac{\partial}{\partial x_{\mu}} \equiv \partial^{\mu} \equiv (\frac{\partial}{\partial t}; -\vec{\nabla})$$

i. Construct the Klein-Gordon operator using the identifications

$$E \to i \frac{\partial}{\partial t} \qquad \qquad \vec{p} \to -i \vec{\nabla}$$

ii. Show that if we define the four-vector current  $J^{\mu} \equiv (\rho; \vec{j})$ , where  $\rho$  and  $\vec{j}$  the charge density and current respectively, the continuity equation in electrodynamics can be written as the conservation of the four-current as

$$\partial_{\mu}J^{\mu} = 0$$

## 2. About those Delta Functions

Using the defining property of the Delta function

$$\int_{-\infty}^{+\infty} f(x)\,\delta(x)\,dx = f(0)$$

for some function f(x), and that  $\delta(-x) = \delta(x)$ 

(a) Show that

$$\int_{-\infty}^{+\infty} f(x)\delta(a\,x)\,dx = \frac{1}{|a|}\,f(0)$$

for an arbitrary constant a.

(b) Using all these show that for a function g(x)

$$\int_{-\infty}^{+\infty} f(x)\,\delta(g(x))\,dx = \sum_{a} \frac{1}{|g'(a)|} \int_{-\infty}^{+\infty} f(x)\,\delta(x-a)\,dx$$

where a are the zeroes of g(x) (i.e.  $g(a) = 0, \forall a$ ) while  $g'(a) \neq 0$ .

(c) Now verify that for  $\omega_p = \sqrt{(\vec{p})^2 + m^2} > 0$  we have

$$\int d^4p \,\delta(P^2 - m^2) \,f(P) = \int d^3p \int dp_0 \,\delta(p_0^2 - (\vec{p})^2 - m^2) \,f(P) = \frac{1}{2\omega_p} \,\int d^3p \,f(\omega_p, \vec{p}) \,d^3p \,f(\omega_p,$$

where you need to use the fact that since the four-momentum is always timelike, the sign of  $p_0$  is Lorentz invariant. Thus we need only integrate  $p_0$  from 0 to  $\infty$ , which in turn means we only need the positive root.

## 3. Electromagnetism:

The covariant form of the Lagrangian density for the electromagnetic field in the absence of sources is

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} ,$$

where  $F_{\mu\nu} = \partial_{\mu}A_{\nu} - \partial_{\nu}A_{\mu}$  is the field strength tensor defined in terms of the potential 4-vector  $A_{\mu}$ .

- (a) Derive the equations of motion for  $A_{\mu}(x)$  assuming it is the dynamical field, using Euler-Lagrange. Show that these are equivalent to Maxwell's equations in the vacuum. (It is useful to remember that  $E^i = -F^{0i}$  and  $\epsilon^{ijk}B^k = -F^{ij}$ .)
- (b) If we now consider sources in the form of a current  $j_{\mu}$ , what is the form of the interaction term between  $A_{\mu}$  and  $j_{\mu}$  we must add to  $\mathcal{L}$  in order to obtain Maxwell's equations in the presence of the source  $j_{\mu} = (\rho, \vec{j})$ ?
- (c) The Lagrangian in the absence of sources is invariant under the gauge transformation

$$A_{\mu}(x) \to A_{\mu}(x) + \partial_{\mu}\alpha(x)$$

with  $\alpha(x)$  an arbitrary function. What is the condition that  $j_{\mu}$  must fulfill for the Lagrangian to remain gauge invariant in the presence of sources ?