

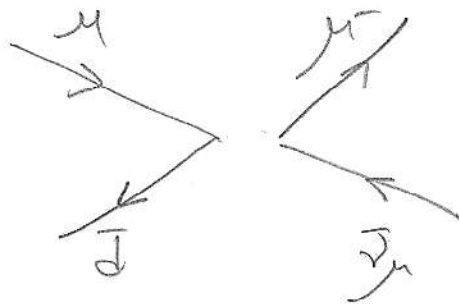
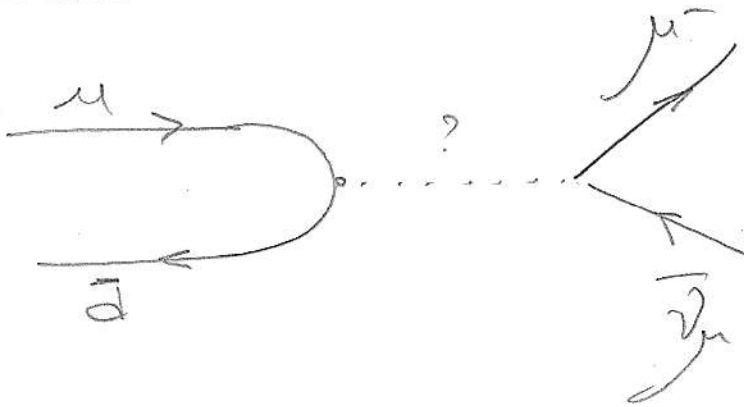
Interações Fracas II

A19 ✓ ⊙

Decaimento do π^-

$$\pi^- \rightarrow \mu^- \bar{\nu}_\mu; e \bar{\nu}_e$$

π^- é $u \bar{d}$!



$$\pi^-(q) \rightarrow \mu^-(p) \bar{\nu}_\mu(k)$$

$$q_\mu = p_\mu + k_\mu$$

$$\mathcal{M} = \frac{G_F}{\sqrt{2}} \bar{d} \gamma_\mu (1 - \gamma_5) u \bar{\mu} \gamma^\mu (1 - \gamma_5) \nu_\mu$$

↳ Processo é um tipo "particular" (de quarks)!

O modo como processo envolve π^- !

→ Escolher a corrente V de quatro ^{componentes} forma covariante

$$j_\mu = \bar{\psi} \gamma_\mu (1 - \gamma_5) \psi \quad \text{tal que}$$

$$\langle \pi^-(q) | j_\mu | 0 \rangle$$

é o elemento de matriz do operador j_μ entre $\langle \pi^-(q) |$ e $| 0 \rangle$! $\pi \rightarrow \text{circulo} \rightarrow \gamma_\mu \pi$

Parametrizar: Não sabemos calcular o elemento de matriz "hadronico". Mas ela deve ter a forma:

$$\langle \pi^-(q) | j_\mu | 0 \rangle = f_\pi * f(q^2)$$

q^2 é um escalar de Lorentz formado com j_μ . MAS $q^2 = m_\pi^2$!

→ Chamamos de $f(q^2 = m_\pi^2) \equiv f_\pi$ constante de desintegração do píon

$$m = \frac{G_F}{\sqrt{2}} f_\pi (p^\mu + k^\mu) \bar{\psi} \gamma_\mu (1 - \gamma^5) \psi$$

Se $m_D \approx 0 \Rightarrow$

$$k^\mu \gamma_\mu \psi = \cancel{k} \psi = -m_D \psi = 0$$

$$\bar{\psi} \cancel{p} \gamma_\mu = \bar{\psi} \cancel{p} = m_D \bar{\psi} \quad \text{vinda}$$

de

$$(\cancel{p} - m) \psi(\vec{p}) = 0$$

$$\bar{\psi}(\vec{p})(\cancel{p} - m) = 0$$

$$\Rightarrow m = \frac{G_F}{\sqrt{2}} f_\pi m_D \bar{\psi} (1 - \gamma^5) \psi$$

$$|m|^2 = \frac{G_F^2}{2} f_\pi^2 m_D^2 \text{Tr} [\cancel{p}(1 - \gamma^5) \cancel{k}(1 + \gamma^5)]$$

$\cancel{p} + m_D$
 ↳ não contrib.

$$= \dots \left\{ \text{Tr}[\cancel{p}\cancel{k}] - \text{Tr}[\cancel{p}\gamma^5\cancel{k}\gamma^5] \right. \\ \left. - \text{Tr}[\cancel{p}\gamma^5\cancel{k}] + \text{Tr}[\cancel{p}\cancel{k}\gamma^5] \right\}$$

Leitender

(4)

$$|M|^2 \sim \bar{\mu}_\alpha (1-\gamma^5)_{\alpha\beta} \nu_\beta \left(\mu^\dagger \gamma^0 (1-\gamma^5) \nu \right)^\dagger$$

$$\sim \bar{\mu}_\alpha (1-\gamma^5)_{\alpha\beta} \nu_\beta \nu^\dagger (1-\gamma^5) \gamma^0 \mu$$

$$\sim \bar{\mu}_\alpha (1-\gamma^5)_{\alpha\beta} \nu_\beta \nu^\dagger \gamma^0 (1+\gamma^5) \mu = \bar{\mu}_\alpha (1-\gamma^5)_{\alpha\beta} \nu_\beta \underbrace{\nu^\dagger (1+\gamma^5)}_{\gamma} \mu = \bar{\mu}_\alpha (1-\gamma^5)_{\alpha\beta} \underbrace{\nu_\beta \gamma}_{\gamma} \mu$$

$$\sim \bar{\mu}_\alpha (1-\gamma^5)_{\alpha\beta} \nu_\beta \gamma (1+\gamma^5) \mu$$

$$\Rightarrow |M|^2 = \frac{G_F^2}{2} f_\pi^2 m_\mu^2 \text{Tr}[\gamma(1-\gamma^5) \gamma (1+\gamma^5)]$$

$$= \frac{G_F^2}{2} f_\pi^2 m_\mu^2 2 \text{Tr}[\gamma \gamma]$$

$$\boxed{|M|^2 = 4 G_F^2 f_\pi^2 m_\mu^2 P \cdot K} = \frac{|M|^2}{\downarrow!}$$

$$\Gamma = \frac{1}{2E_{\pi}} |\overline{M}|^2 \frac{d^3 p_{\mu}}{(2\pi)^3 2E_{\mu}} \frac{d^3 p_{\nu}}{(2\pi)^3 2E_{\nu}} (2\pi)^4 \delta^{(4)}(P_A - P_{\mu} - P_{\nu}) \quad (5)$$

Aqui E_{π} pode ser no sistema em repouso
 $E_{\pi} = m_{\pi}$

$$\Gamma = \frac{1}{2 \cdot (2\pi)^2 m_{\pi}} |\overline{M}|^2 \frac{d^3 p_{\mu}}{E_{\mu} E_{\nu}} \delta(E_{\pi} - E_{\nu} - E_{\mu})$$

$$\Gamma = \frac{1}{8\pi^2 m_{\pi}} |\overline{M}|^2 \frac{p_{\mu}^2 d p_{\mu}}{E_{\mu} E_{\nu}} \sqrt{4\pi} \delta(m_{\pi} - E_{\nu} - E_{\mu})$$

Para fazer o integral:

$$\delta(m_{\pi} - E_{\nu} - E_{\mu}) = \delta(m_{\pi} - \sqrt{p_{\mu}^2 + m_{\nu}^2} - \sqrt{p_{\mu}^2 + m_{\nu}^2})$$

$$\Rightarrow \Gamma = \frac{4\pi}{8\pi^2 m_{\pi}} |\overline{M}|^2 \frac{p_{\mu}^2}{E_{\mu} E_{\nu}} \frac{1}{\left(\frac{p_{\mu}}{E_{\mu}} + \frac{p_{\mu}}{E_{\nu}}\right)}$$

$$\Gamma = \frac{1}{2\pi m_{\pi}} |\overline{M}|^2 \frac{p_{\mu}}{[E_{\mu} + E_{\nu}]}$$

$= m_{\pi}$

$$\Gamma = \frac{1}{2\pi m_{\pi}^2} p_{\mu} |\overline{m}|^2$$

onde $p_{\mu} \equiv m_{\pi} = E_{\nu} + E_{\mu}$

$$m_{\pi} = p_{\mu} + \sqrt{p_{\mu}^2 + m_{\mu}^2}$$

$$m_{\pi} - p_{\mu} = \sqrt{p_{\mu}^2 + m_{\mu}^2}$$

$$m_{\pi}^2 + p_{\mu}^2 - 2m_{\pi}p_{\mu} = p_{\mu}^2 + m_{\mu}^2$$

$$\Rightarrow p_{\mu} = \frac{m_{\pi}^2 - m_{\mu}^2}{2m_{\pi}}$$

$$\Rightarrow \Gamma = \frac{(m_{\pi}^2 - m_{\mu}^2)}{4\pi m_{\pi}^3} |\overline{m}|^2$$

$$|\overline{m}|^2 = 4G_F^2 f_{\pi}^2 m_{\mu}^2 P \cdot K =$$

$$P \cdot K = E_{\mu} E_{\nu} - \vec{p}_{\mu} \cdot \vec{p}_{\nu} = E_{\mu} E_{\nu} + p_{\mu}^2$$

=

$$P_{\alpha} k = E_{\mu} P_{\mu} + P_{\mu}^2$$

(7)

$$E_{\mu} = \sqrt{P_{\mu}^2 + m_{\mu}^2}$$

$$E_{\mu}^2 = \frac{(m_{\pi}^2 - m_{\mu}^2)^2}{4m_{\pi}^2} + \frac{4m_{\pi}^2 m_{\mu}^2}{4m_{\pi}^2}$$

$$E_{\mu}^2 = \frac{(m_{\pi}^2 + m_{\mu}^2)^2}{4m_{\pi}^2} \Rightarrow E_{\mu} = \frac{(m_{\pi}^2 + m_{\mu}^2)}{2m_{\pi}}$$

$$\Rightarrow P_{\alpha} k^{\alpha} = \frac{(m_{\pi}^2 + m_{\mu}^2)}{2m_{\pi}} \frac{(m_{\pi}^2 - m_{\mu}^2)}{2m_{\pi}} + \frac{(m_{\pi}^2 - m_{\mu}^2)^2}{4m_{\pi}^2}$$

$$= \frac{(m_{\pi}^2 - m_{\mu}^2)}{4m_{\pi}^2} \left[m_{\pi}^2 + m_{\mu}^2 + m_{\pi}^2 - m_{\mu}^2 \right]$$

$$P_{\alpha} k^{\alpha} = \frac{(m_{\pi}^2 - m_{\mu}^2)}{2}$$

$$\Gamma = \frac{(m_{\pi}^2 - m_{\mu}^2)^2}{8\pi m_{\pi}^3} 4 G_F^2 f_{\pi}^2 m_{\mu}^2$$

⊗

$$\Gamma = \frac{G_F^2}{2\pi} f_{\pi}^2 m_{\mu}^2 m_{\pi} \left(1 - \frac{m_{\mu}^2}{m_{\pi}^2}\right)^2$$

$f_{\pi} = ?$ Medida experimentalmente de aqui!

$\sim f_{\pi} \sim 132 \text{ MeV} \sim m_{\pi}!$

$$G_F = 1.166 \cdot 10^{-5} \text{ GeV}^{-2}$$

$$\Rightarrow \Gamma = 5.28 \cdot 10^{-16} \left(1 - \frac{m_{\mu}^2}{m_{\pi}^2}\right)^2 \text{ GeV}$$

$$\Gamma = 1.27 \cdot 10^{-16} \text{ GeV}$$

$$\tau = \frac{\hbar}{\Gamma} = \frac{6.582 \cdot 10^{-25} \text{ GeV s}}{1.27 \cdot 10^{-16} \text{ GeV}}$$

$$\tau_{\pi} \approx 5 \times 10^{-9} \text{ s} \quad \text{A verdade } \sim 2.8 \cdot 10^{-8} \text{ s} \quad ??$$

Supressão de Γ por helicidade:

(9)

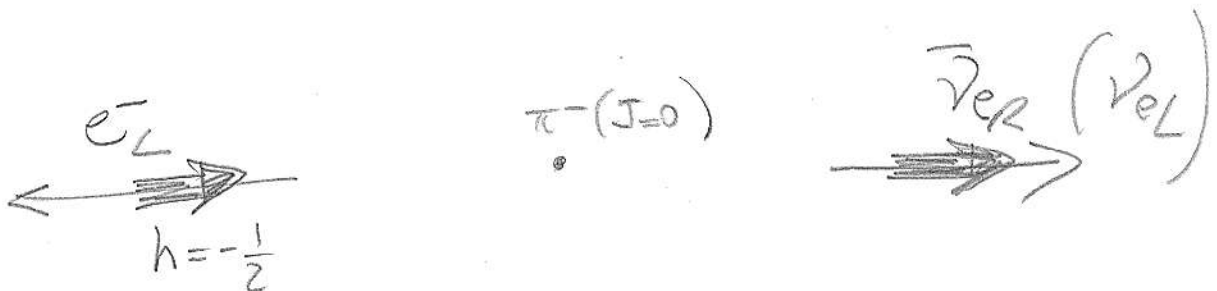
$$\frac{\Gamma(\pi^- \rightarrow \mu^- \bar{\nu}_\mu)}{\Gamma(\pi^- \rightarrow e^- \bar{\nu}_e)} = \frac{m_\mu^2}{m_e^2} \frac{(m_\pi^2 - m_\mu^2)^2}{(m_\pi^2 - m_e^2)^2}$$

$$\Rightarrow \Gamma(\pi^- \rightarrow \mu^- \bar{\nu}_\mu) \approx 10^4 \Gamma(\pi^- \rightarrow e^- \bar{\nu}_e)!$$

\Rightarrow π^- decai em μ^- ! quase sempre!

Por que!? No limite $m_e \rightarrow 0$

π^- Não decai!



Se o e não decai para π^0 -esquerda ($\bar{\nu}_e$ é de π^0 -direita)

\Rightarrow Não pode decair por cons. de momento!

$$J_\pi = 0 \quad J_{\text{final}} = 1 \quad ! \quad ? \quad ?$$

Pneira com r!

(10)



MAS as interações são de acoplam com e_L !
 \Rightarrow Como fazer para $e_L \rightarrow e_R$?

Termo de MASSA!

$$\mathcal{L} = \bar{e}_L i \not{\partial} e_L + \bar{e}_R i \not{\partial} e_R - m_e (\bar{e}_L e_R + \bar{e}_R e_L)$$



Por isso que no limite $m_e \rightarrow 0$ (ou $m_e \rightarrow \infty$)
Não temos decaimento do π^- !

Esta característica do π^- é inerente às
interações fracas!